

An explicit formula of powers of the 2 × 2 quantum matrices and its applications

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MSC classes:16T20, 33C45, 81R50

Abstract

We present an explicit formula of the powers for the 2×2 quantum matrices, that is a natural quantum analogue of the powers of the usual 2×2 matrices. As applications, we give some non-commutative relations of the entries of the powers for the 2×2 quantum matrices, which is a simple proof of the results of Vokos-Zumino-Wess (1990).

1 Introduction

Let *A* be a 2×2 matrix over a fixed base field *k* and *a*,*b*,*c*,*d* \in *k* are entries of *A*:

$$A = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

For any positive number *n*, the following explicit formula of *A*^{*n*} holds:

$$A^{n} = \begin{pmatrix} af_{n-1}(\operatorname{tr} A) - \det Af_{n-2}(\operatorname{tr} A) & bf_{n-1}(\operatorname{tr} A) \\ cf_{n-1}(\operatorname{tr} A) & df_{n-1}(\operatorname{tr} A) - \det Af_{n-2}(\operatorname{tr} A) \end{pmatrix}$$
(1.1)

¹ Dedicated to T. Umeda and M. Wakayama for their 66th birthdays.

Here $f_n(x)$ is the polynomial of degree *n* defined by

$$f_n(x) = f_n(x,y) := \sum_{l=0}^{\lfloor \frac{n}{2} \rfloor} (-1)^l \binom{n-l}{l} x^{n-2l} y^l, \quad f_{-1}(x) := 0$$
, (1.2) which is

the Chebyshev polynomial of the second kind:

$$U_n(x) := \sum_{l=0}^{\lfloor \frac{n}{2} \rfloor} (-1)^l \binom{n-l}{l} x^{n-2l}$$

In fact this formula is well-known in linear algebra, and its proof is easy by induction and the recurrence relation for $f_n(x)$:

$$f_{n+1}(x) = x f_n(x) - y f_{n-1}(x).$$
(1.3)

In this note, we give a quantum analogue of the formula (1.1) and its applications. First, we review some fundamental objects and facts on quantum matrix or groups [M], [T] relevant to the main results of this paper.

We call *A* a 2 × 2 (*q*-)quantum matrix if its entries satisfy the following relations:

$$ab = qba$$
, $ac = qca$, $ad - da = (q - q^{-1})bc$,
 $bc = cb$, $bd = qdb$, $cd = qdc$, (R_q)

Here q is a central indeterminate. A quantum analogue of the coordinate ring $A_q(Mat(2))$ is the algebra generated by a,b,c,d and q which is a typical example of quantum groups.

The quantum adjoint matrix of any quantum matrix A is defined as

$$\hat{A} = \begin{pmatrix} \hat{a} & \hat{b} \\ \hat{c} & \hat{d} \end{pmatrix} := \begin{pmatrix} d & -q^{-1}b \\ -qc & a \end{pmatrix}$$
(1.4)

By the definition of A^{\wedge} and the relations (R_q), the quantum adjoint matrix A^{\wedge} satisfies the relations (R_{q-1}):

$$a^{h}b = q^{-1}ba, a^{h}c^{h} = q^{-1}c^{h}a, a^{h}d^{h} - d^{h}a^{h} = (q^{-1}-q)^{h}bc,$$

 $bc^{h}c^{h} = c^{h}b, bd^{h} = q^{-1}d^{h}b, c^{h}d^{h} = q^{-1}d^{h}c.$

Hence, the relations (R_q) are equivalent to

$$bc = cb, \quad A\hat{A} = \hat{A}A = \begin{pmatrix} \delta & 0\\ 0 & \delta \end{pmatrix} = \delta E_2$$
(1.5)

where E_2 is the 2×2 identity matrix, and $\delta := ad-qbc = da-q^{-1}bc$ is the quantum determinant of *A* which is a central element of $A_q(Mat(2))$.

For convenience, we introduce a 2×2 matrix

$$C := \begin{pmatrix} q^{\frac{1}{2}} & 0\\ 0 & q^{-\frac{1}{2}} \end{bmatrix}$$

and put

$$\tau := \operatorname{tr} (AC) = q^{\frac{1}{2}}a + q^{-\frac{1}{2}}d, \quad \tau' := \operatorname{tr} (C^{-1}A) = q^{-\frac{1}{2}}a + q^{\frac{1}{2}}d.$$

Our main results are following.

Theorem 1.1. For any positive integer *n*, we have

$$A_{n} = AC_{-n+1}f_{n-1}(\tau) - C_{-n}\delta f_{n-2}(\tau)$$
(1.6)

$$= f_{n-1}(\tau')C_{n-1}A - f_{n-2}(\tau')\delta C_{n}, \qquad (1.7)$$

where

$$f_n(\tau) := f_n(\tau, \delta) = \sum_{l=0}^{\lfloor \frac{n}{2} \rfloor} (-1)^l \binom{n-l}{l} \tau^{n-2l} \delta^l, \quad f_{-1}(\tau) := 0$$

Let us put

$$A^{n} = \begin{pmatrix} a_{n} & b_{n} \\ c_{n} & d_{n} \end{pmatrix}, \quad \hat{A}^{m} := \begin{pmatrix} \hat{a}_{m} & \hat{b}_{m} \\ \hat{c}_{m} & \hat{d}_{m} \end{pmatrix}.$$

By comparing the entries of A^n and (1.6), (1.7), we obtain the following explicit formulas of the entries of A^n .

Corollary 1.2. For any positive integer *n*, we have

$$a_{n} = q^{-} \underline{}^{2} a f_{n-1}(\tau) - q^{-} \underline{}^{2} \delta f_{n-2}(\tau) = q \underline{}^{2} f_{n-1}(\tau') a - q \underline{}^{2} \delta f_{n-2}(\tau'),$$
(1.8)

$$b_n = q \underline{\ }_{n-1}^2 b f_{n-1}(\tau) = q \underline{\ }_{n-1}^2 f_{n-1}(\tau') b, \tag{1.9}$$

$$c_n = q_{-2}c_{f_{n-1}}(\tau) = q_{-2}c_{f_{n-1}}(\tau')c,$$
 (1.10)

$$n - 1 n n n - 1 n dn = q __2 df n - 1(\tau) - q __2 \delta f n - 2(\tau) = q - __2 f n - 1(\tau') d - q - __2 \delta f n - 2(\tau').$$
(1.11)

As applications, we give the following results, in particular Theorem 1.4. **Corollary 1.3.** For any positive integer *m*, we have

$$a_{m}^{*} = \frac{m-1}{q^{2} a f} q^{2} a f \frac{m}{m-1} (\tau) - q^{2} \delta f_{m-2}(\tau)$$

$$= q^{-\frac{m-1}{2}} f_{m-1}(\tau') \hat{a} - q^{-\frac{m}{2}} \delta f_{m-2}(\tau') = d_{m}, \qquad (1.12)$$

$$\hat{b}_{m} = q^{-\frac{m-1}{2}} \hat{b}_{m-1}(\tau) = q^{-\frac{m-1}{2}} f_{m-1}(\tau') \hat{b} = -q^{-m} b_{m}, \qquad (1.13)$$

$$c_{m}^{*} = q \underbrace{-}_{2}^{2} cf_{m-1}(\tau) = q \underbrace{-}_{2}^{2} f_{m-1}(\tau') c^{*} = -q^{m} c_{m}, d^{*} m = q \underbrace{-}_{m^{2}-1} df_{m-1}(\tau) - q \underbrace{-}_{m^{2}} \delta f_{m-2}(\tau) = q^{m} \underbrace{-}_{2}^{-} f_{m-1}(\tau') d^{*} - q^{-\frac{m}{2}} \delta f_{m-2}(\tau') = a_{m}.$$
(1.14)

Theorem 1.4 ([VZW]). For any non-negative integers *m* and *n*, we have

 $\delta^m a$

$$dman - q - mbmcn = andm - q_mbncm = (\delta n dm_n - nm \quad ((m < nm \ge n)), \quad (1.16)$$

$$\delta^{m}b$$

$$(dmbn - q - mbmdn = -q - manbm + bnam = -nn - mm\delta nbm - n \quad ((m < nm \ge n)),$$

$$q$$

$$(1.17)$$

$$m$$

$$(m < n)$$

$$-q \ cman + amcn = cndm - q \ dncm^{=} - qn - m\delta n_{Cm^{-}n} \qquad (m \ge n),$$
(1.18)

$${}^{m} - {}^{m}\delta^{m}d_{n-m} \quad (m < n)$$

$$-q \operatorname{CmDn} + aman = -q \operatorname{CnDm} + anam \quad Onam^{-n} \qquad (m \ge n), \quad (1.19)$$

 $b_n c_m - q^{n-m} c_n b_m = 0.$

2 **Proof of Theorem 1.1**

To prove Theorem 1.1, we need a quantum analogue of Cayley-Hamilton theorem.

Lemma 2.1 ([UW] Lemma 3). The following formula holds.

$$A^{2} = AC^{-1}\tau - C^{-2}\delta = \tau'CA - \delta C^{2}.$$
 (2.1)

Proof of Theorem 1.1 Since (1.6) and (1.7) can be similarly proved, we only prove (1.6). These formulas are proved by induction on *n*.

The n = 1 case is trivial. Assume the case of n holds. Hence, from the induction hypothesis we have

$$A^{n+1} = AA^{n}$$

= $A \left\{ AC^{-n+1} f_{n-1}(\tau) - C^{-n} \delta f_{n-2}(\tau) \right\}$
= $A^{2}C^{-n+1} f_{n-1}(\tau) - AC^{-n} \delta f_{n-2}(\tau).$

By Cayley-Hamilton theorem (2.1) and the recursion (1.3), we have

$$A_{n+1} = (AC_{-1}\tau - C_{-2}\delta)C_{-n+1}f_{n-1}(\tau) - AC_{-n}\delta f_{n-2}(\tau)$$

= $AC_{-n}(\tau f_{n-1}(\tau) - \delta f_{n-2}(\tau)) - C_{-n-1}\delta f_{n-1}(\tau) = AC_{-n}f_n(\tau)$
- $C_{-n-1}\delta f_{n-1}(\tau).$

The formula (1.7) can be proved by the similar argument for $A^{n+1} = A^n A$.

(1.20)

Corollary 1.2 follows from comparing the entries of A^n and (1.6), (1.7) immediately.

Remark 2.2. (1) By consider the classical limit q = 1 in Corollary 1.2, we recover the classical result (1.1).

(2) From the recursion (1.3), we derive other expressions of a_n and d_n :

$$a_{n} = q^{-2} f_{n}(\tau) - q^{-2} df_{n-1}(\tau) = q^{-2} f_{n}(\tau') - q^{2} f_{n-1}(\tau') d, \qquad (2.2)$$

$$d_n = q_{-2}f_n(\tau) - q^2 a f_{n-1}(\tau) = q^{-2}f_n(\tau') - q^{-2}f_{n-1}(\tau')a.$$
(2.3)

Since these expressions (2.2), (2.3) (and (1.9), (1.10)) hold for n = 0, Corollary 1.2 is also true for the case of n = 0. (3) Umeda-Wakayama [UW] considered

$$\tau_{n} := \operatorname{tr} A C = q_{2} a_{n} + q_{2} b_{n},$$

$$\tau_{n'} := \operatorname{tr} C - nA_{n} = q_{-n^{2}} a_{n} + q_{-n^{2}} b_{n},$$

and pointed out that τ_n and τ_n' satisfy the following Fibonacci type equations:

$$\tau_{n+1} = \tau_n \tau - \tau_{n-1} \delta, \quad \tau'_{n+1} = \tau'_n \tau' - \tau'_{n-1} \delta.$$
 (2.4)

These equations (2.4) are equal to the recursion (1.3) of $f_n(\tau)$ exactly. Hence by $\tau_1 = \tau$ and $\tau'_1 = \tau'$ we have

$$\tau_n = f_n(\tau), \quad \tau'_n = f_n(\tau')$$
(2.5)

3 Applications

We point out that quantum adjoint matrix A^{-1} is a q^{-1} -quantum matrix and

$$\tau^{\hat{}} := q_{-21}a^{\hat{}} + q_{12}d^{\hat{}} = q_{12}a + q_{-21}d = \tau \tau^{\hat{}}$$
$$:= q_{21}a^{\hat{}} + q_{-21}d^{\hat{}} = q_{-12}a + q_{21}d = \tau', \ \delta^{\hat{}} := a^{\hat{}}d^{\hat{}} - q^{-1}bc^{\hat{}} = da - q^{-1}bc = \delta.$$

Then we prove Corollary 1.3.

From (1.5), Corollary 1.2 and Corollary 1.3, we obtain the proof of Theorem 1.4 which is a simple prood of Vokos-Zumino-Wess [VZW].

Proof of Theorem 1.4 For any non-negative integers *m*,*n*, from Corollary 1.3 we have

$$\hat{A}^{m}A^{n} = \begin{pmatrix} d_{m} & -q^{-m}b_{m} \\ -q^{m}c_{m} & a_{m} \end{pmatrix} \begin{pmatrix} a_{n} & b_{n} \\ c_{n} & d_{n} \end{pmatrix}$$
$$= \begin{pmatrix} d_{m}a_{n} - q^{-m}b_{m}c_{n} & d_{m}b_{n} - q^{-m}b_{m}d_{n} \\ -q^{m}c_{m}a_{n} + a_{m}c_{n} & -q^{m}c_{m}b_{n} + a_{m}d_{n} \end{pmatrix}$$
(3.1)

and

$$A^{n}\hat{A}^{m} = \begin{pmatrix} a_{n} & b_{n} \\ c_{n} & d_{n} \end{pmatrix} \begin{pmatrix} d_{m} & -q^{-m}b_{m} \\ -q^{m}c_{m} & a_{m} \end{pmatrix}$$
$$= \begin{pmatrix} a_{n}d_{m} - q^{m}b_{n}c_{m} & -q^{-m}a_{n}b_{m} + b_{n}a_{m} \\ c_{n}d_{m} - q^{m}d_{n}c_{m} & -q^{-m}c_{n}b_{m} + d_{n}a_{m} \end{pmatrix}.$$
(3.2)

On the other hand, by applying $AA^{\hat{}} = AA^{\hat{}} = \delta E_2$ we obtain

 $A^{M} = AnA^{M}$

$$\boxed{222} \boxed{2} \delta m A_{n-m} = \delta m C_{n-n-mm} - \delta \delta m^{-m} db_{nn-mm}! - - -$$

$$-!(m < n) \delta^m a$$

$$= .$$

$$nA^{n}m = -qm\delta ndnm\delta nCnm n - qn\delta nmam\delta nbnm n \qquad (m \ge n) \delta \ \text{PP}$$

$$(3.3)$$

By comparing the entries of (3.1), (3.2) and (3.3) we have (1.16), (1.17), (1.18) and (1.19).

Finally, the relation (1.20) follows from the explicit formulas (1.9) and (1.10):

$$bncm = \frac{\frac{n-1}{q}}{q} \frac{\frac{m-1}{2}}{2} f_{n-1}(\tau)bq - 2cfm - 1(\tau)$$
$$= q q q \frac{2f_{n-1}(\tau)cq 2bf_{m-1}(\tau)}{2}$$
$$= qn-mcnbm.$$

If we set m = n in (1.16), (1.17), (1.18), (1.19) and (1.20), then we obtain an interesting Corollary which means that A^n is a 2×2 q^n -quantum matrix.

Corollary 3.1. For any non-negative integer *n*, we have *anbn* = *qnbnan*, *ancn*

 $= q_n c_n a_n, a_n d_n - d_n a_n = (q_n - q_{-n}) b_n c_n,$ $b_n c_n = c_n b_n, b_n d_n = q_n d_n b_n, c_n d_n = q_n d_n c_n \text{ and} \qquad (R_{q_n})$

(3.4)

 $andn - qnbncn = dnan - q-nbncn = \delta n$. i.e. (the quantum determinant of A^n) = (the quantum determinant of A)^{*n*}

Originally, Theorem 1.4 was proved by Vokos-Zumino-Wess [VZW] and its proof was a brute force approach using double induction on m and n. Later, Corrigan-Fairlie-Fletcher-Sasaki [CFFS] and Umeda-Wakayama [UW] gave some simple proofs of Corollary 3.1 which is the case of m = n in Theorem 1.4 independently. Our proof of Theorem 1.4 is different from any of them.

In particular, it is desirable to extend Theorem 1.1 and Corollary 1.2 to $n \times n$ quantum matrices.

Acknowledgement

We would like to thank Professor To^ru Umeda (Osaka City University) and Professor Masato Wakayama (Nippon Telegraph and Telephone Corporation, Institute for Fundamental Mathematics) for their valuable comments for references and various techniques of quantum calculus. This work was supported by Grant-in-Aid for Young Scientists (Number 21K13808).

References

- [CFFS] E.Corrigan, D.B.Fairlie, P.Fletcher and R.Sasaki : *Some aspects of quantum groups and supergroups*, J. Math. Phys. **31** (1990), 776–780.
- [M] Y.I.Manin: *Topics in noncommutative geometry*, Princeton Univ. Press, 1991.
- [T] M.Takeuchi: A Short Course on Quantum Matrices, New Directions in Hopf Algebras, MSRI Publications 43, Cambridge University Press (2002), 383-435.

- [UW] T.Umeda and M.Wakayama : *Powers of* 2 × 2 *quantum matrices*, Comm. Alg. **21** (1993), 4461–4465.
- [VZW] S.P.Vokos, B.Zumino and J.Wess : *Analysis of the basic matrix representation of* $GL_q(2,C)$, Z. Phys. C **48** (1990), 65-74.